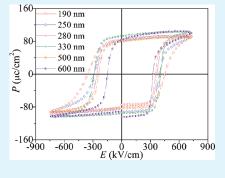
# Ferroelectric Behavior in Bismuth Ferrite Thin Films of Different Thickness

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ABSTRACT: The ferroelectric behavior of BiFeO<sub>3</sub> thin films is modified by changing the film thicknesses, where the BiFeO<sub>3</sub> thin films with different thicknesses were grown on SrRuO<sub>3</sub>/Pt/TiO<sub>2</sub>/SiO<sub>2</sub>/Si(100) substrates by radio frequency sputtering. The mixture of (110) and (111) orientations is induced for all BiFeO<sub>3</sub> thin films regardless of their thicknesses, together with the columnar structure and the dense microstructure. Their dielectric behavior is almost independent of the film thickness where all thin films have a low dielectric loss. A giant remanent polarization of  $2P_r \approx 156.6-188.8 \,\mu C/cm^2$  is induced for the BiFeO<sub>3</sub> thin films in the thickness range of 190–600 nm. As a result, it is an effective way to improve the ferroelectric behavior of the BiFeO<sub>3</sub> thin film by tailoring the film thickness.



KEYWORDS: bismuth ferrite, thickness dependence, ferroelectric properties

# **1. INTRODUCTION**

Considerable attention has been given to multiferroic (BiMnO<sub>3</sub>, TbMnO<sub>3</sub>, and BiFeO<sub>3</sub>) materials because of the coexistence of ferroelectricity, (anti)ferromagnetism, and ferroelasticity.<sup>1-6</sup> Recently, the BiFeO<sub>3</sub> (BFO) material has been extensively studied from the viewpoint of device applications,  $5^{-14}$  because of the Neel temperature ( $T_{\rm N} \approx 370$  °C), a high Curie temperature ( $T_{\rm c} \approx$ 850 °C), and a giant polarization value confirmed by the theories<sup>15</sup> and the experiments.<sup>16,17</sup> According to the theoretical predication, an unusual spontaneous polarization (up to  $\sim 100 \ \mu C/cm^2$ ) has been predicted for the BFO material with (111) orientation.<sup>15,16</sup> A larger polarization has also been demonstrated for the epitaxial BFO thin films grown on SrRuO<sub>3</sub>/SrTiO<sub>3</sub> substrates by pulse laser deposition,<sup>6,16</sup> but it is not integrated with other silicon devices because of the introduction of an expensive SrTiO<sub>3</sub> substrate. Moreover, a polycrystalline BFO thin film with a tetragonal structure is of a giant polarization value of  $P_{\rm r} \approx 150 \ \mu {\rm C/cm}^2$ when prepared on the  $Pt/TiO_2/SiO_2/Si(100)$  substrate by pulsed laser deposition, but the measurement temperature is below room temperature (90 K).<sup>13</sup> Therefore, the expensive SrTiO<sub>3</sub> substrate and the low-temperature measurement technique hinder the practical applications considered for the bismuth ferrite thin films.

Some attempts have been conducted to improve the electrical properties of bismuth ferrite thin films, such as the employment of the film thickness effects,<sup>18–22</sup> ion substitution,<sup>14</sup> buffer layer,<sup>23</sup> and orientation.<sup>16,17</sup> There have been some works about the thickness dependence of electrical properties in BFO thin films by tailoring the film thicknesses, as listed in Table 1.<sup>6,18–22</sup> A high remanent polarization is often demonstrated for the BFO thin film with an optimum film thickness when grown on SrRuO<sub>3</sub>/SrTiO<sub>3</sub> substrates.<sup>6,18,19</sup> However, BFO thin films grown in oxides buffer layer/silicon substrates exhibit a poorer ferroelectric behavior,<sup>20,21</sup> or the ferroelectric properties can be only measured at a low temperature.<sup>22</sup> It is highly desirable to prepare the BFO thin films

on a silicon substrate to enable the integration with other Si devices, and the trend of miniaturization of devices requires the reduction in the thickness of thin films. However, the direct growth of BFO thin films on Pt-coated silicon substrate often results in the formation of a dead layer at the interface and the generation of secondary phases,<sup>23</sup> but BFO thin films can be well grown on Pt-coated silicon via several oxide buffer layers including SrRuO<sub>3</sub>,<sup>17</sup> BaPbO<sub>3</sub>,<sup>20</sup> LaNiO<sub>3</sub>,<sup>24</sup> and (La, Sr)MnO<sub>3</sub>.<sup>25</sup>

In the present work, we modify the thickness of BFO thin films grown on oxide buffer layer  $SrRuO_3/Pt/TiO_2/SiO_2/Si(100)$  substrates by radio frequency (rf) sputtering. The thickness dependence of ferroelectric behavior in BFO thin films is studied. A giant polarization is demonstrated for the BFO thin films grown on silicon substrates by modifying the film thickness.

### 2. EXPERIMENTAL PROCEDURE

A ceramic target of Bi<sub>1.15</sub>FeO<sub>3</sub> was sintered by the mixed oxides of Bi<sub>2</sub>O<sub>3</sub> and Fe<sub>2</sub>O<sub>3</sub>. These powders were mixed by the ball milling for 24 h, and then the powder was calcined at ~720 °C for 6 h after drying. The calcined powder mixture was pressed into a two-inch pellet and then the pellet was sintered at ~820 °C for 2 h. The SrRuO<sub>3</sub> buffer layer with a thickness of ~90 nm is first grown on Pt/TiO<sub>2</sub>/SiO<sub>2</sub>/Si(100) substrates by rf sputtering, and then the BFO thin films with the thickness range of 100–600 nm were deposited on SrRuO<sub>3</sub>/Pt/TiO<sub>2</sub>/SiO<sub>2</sub>/Si(100) substrates by rf sputtering at the substrate temperatures of ~570 °C. They were both deposited under a rf power of 120 W, and at a base pressure of  $3.0 \times 10^{-6}$  Torr and a deposition pressure of 0.20 mm in diameter were prepared on the thin films by rf sputtering through a shadow mask for measuring electrical properties of thin films.

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Table 1. Thick	ness-Dependent Electri	ical Properties for the	Bismuth Ferrite Thin Films
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composition	substrate	thickness (nm)	preparation method	orientation	$2P_{\rm r}$ ( $\mu$ C/cm <sup>2</sup> )	ref
BFO	SrRuO <sub>3</sub> /SrTiO <sub>3</sub> (100)		PLD	(100)	110	6
BFO	SrRuO <sub>3</sub> /SrTiO <sub>3</sub> (001)	450	RF	(001)	170	18
Mn-BFO	SrRuO <sub>3</sub> /SrTiO <sub>3</sub> (001)	70	PLD	(001)	130	19
BFO	BaPbO <sub>3</sub> /Pt/Ti/SiO <sub>x</sub> /Si	230	RF	polycrystal	85.6	20
BFO	LaNiO <sub>3</sub> /Pt/Ti/SiO <sub>2</sub> /Si	115	sol-gel	polycrystal	28.6	21
BFO	Pt/Ti/SiO <sub>2</sub> /Si	400	CSD	polycrystal	180 at 80 K	22
BFO	SrRuO <sub>3</sub> /Pt/Ti/SiO <sub>2</sub> /Si	330	RF	(110)/(111)	$188.8\pm0.2$	in this work

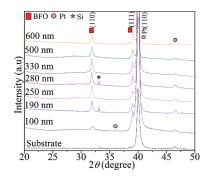


Figure 1. XRD patterns of BiFeO3 thin films with different thicknesses.

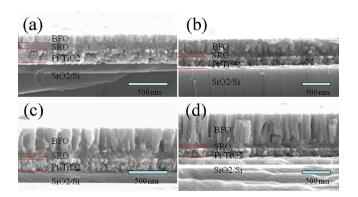


Figure 2. Cross-section of BiFeO<sub>3</sub> thin films with different thicknesses of (a) 100, (b) 250 nm, (c) 330 nm, and (d) 600 nm.

Crystal structures of thin films were measured by X-ray diffraction (XRD) analysis using DX1000 (Dandong, China). Scanning electron microscopy (SEM) (Philips, XL30) was employed to study the cross section of thin films. An impedance analyzer (Solartron Gain phase Analyzer) was employed to measure their dielectric properties, and the ferroelectric behavior is measured by RT2000 Tester (USA).

#### 3. RESULTS AND DISCUSSION

Figure 1 shows the XRD patterns for the BFO thin films with different thicknesses. All thin films possess a pure phase, and no secondary phases are detected. Moreover, the mixture of (110) and (111) orientations is demonstrated for all thin films, as shown by the strong and sharp (110) and (111) peaks, independent of the film thicknesses.

Figure 2 shows the SEM cross sectional images for the BFO thin films with different thicknesses. All thin films have a columnar structure, and the interface among thin film, SRO buffer layer, and Pt electrode is well-established, indicating that

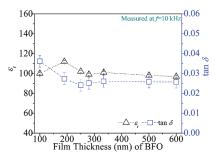
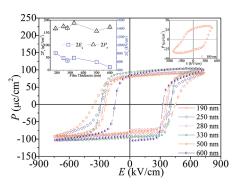


Figure 3. Dielectric behavior of  $BiFeO_3$  thin films with different thicknesses, measured at 10 kHz.

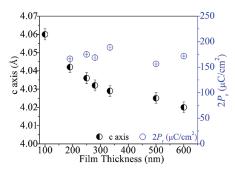


**Figure 4.** P-E loops of BiFeO<sub>3</sub> thin films with different thicknesses, where the inserts are  $2P_r$  and  $2E_c$  values as a function of applied electric fields and the P-E loop for the BFO thin film with a 100 nm thickness.

the BFO layer is well-sputtered on the bottom of the SRO buffer layer. The substantial diffusion between two different phases is not observed. The typical columnar structure in this work confirms that the BFO thin film is well-crystallized.

Figure 3 plots the dielectric constant ( $\varepsilon_r$ ) and the dielectric loss (tan  $\delta$ ) of BFO thin films as a function of the film thicknesses, measured at the frequency of 10 kHz and room temperature. The dielectric behavior slightly changes for all BFO thin films in the range of the film thicknesses investigated, confirming that the dielectric behavior is almost independent of the thicknesses of BFO thin films in the present work. The dead layer forming in the interface between the film and bottom electrode is often responsible for the decrease in  $\varepsilon_r$ . In the present work, the interface between BFO and SRO is well established in Figure 2, resulting in a constant of  $\varepsilon_r$  value. Moreover, the dense microstructure in Figure 2 contributes to the low tan  $\delta$  value in these BFO thin films.

Figure 4 plots the P-E loops for the BFO thin films with different thicknesses, measured at f = 1 kHz and room



**Figure 5.** Summary of the thickness dependence of out-of-plane lattice parameter and polarization.

temperature. All BFO thin films exhibit similar ferroelectric behavior except for the P-E loop of the BFO thin film with a  $\sim$ 100 nm thickness, where a low remanent polarization is induced because of a higher leakage current density of too thin thickness (100 nm) caused by a tunneling effect, as shown in the insert of Figure 4. The insert in Figure 4 shows the  $2P_r$  and  $2E_c$ values as a function of applied electric fields for the BFO thin films with the thicknesses of 190-600 nm. The  $2P_r$  values almost keep unchanged for all thin films with the thicknesses of 190-600 nm. Moreover, the thicker BFO thin films, the smaller  $2E_c$  values. The  $2E_c$  value is  $\sim 808$  kV/cm for the BFO thin film with 190 nm thickness, and then the  $2E_c$  value decreases with an increase in the film thickness. This phenomenon is attributed to the change in the strain relaxation with the film thickness.<sup>19,26,27</sup> In the present work, the  $2P_r$  value is much larger than those of BFO thin films reported by other authors.<sup>6,10,12,14,18–22</sup> It has been reported that the polarization of BFO thin films is almost independent of the strain, as confirmed by the first-principle calculation and the experimental result.<sup>27,28</sup> In this work, the polarization of BFO thin films with 190-600 nm thickness is weakly dependent on the strain, as shown in Figure 5, where the ferroelectric behavior of BFO thin films with 100 nm thickness cannot be measured at room temperature due to the electrical conductivity. The improvement in ferroelectric properties of BFO thin films can be attributed to the mixture of (111) and (110) orientations in this work, where the largest spontaneous polarization happens in the (111) direction for the BFO thin films because that the ions are displaced along [111] directions from their equilibrium positions.<sup>29,30</sup> Therefore, the mixture of (110) and (111) orientations largely contributes to the enhancement in ferroelectric properties of BFO thin films. The columnar structure and dense microstructure also partly contribute to the reduction in the leakage current density of BFO thin films, partly resulting in the improvement in its ferroelectric properties.

In summary, the effect of the thickness on the ferroelectric behavior of BiFeO<sub>3</sub> thin films, which were grown on SrRuO<sub>3</sub>/Pt/TiO<sub>2</sub>/SiO<sub>2</sub>/Si(100) substrates by radio frequency sputtering, is investigated. The column structure and dense microstructure are demonstrated for all BiFeO<sub>3</sub> thin films with (110) and (111) orientations, together with a low dielectric loss. A giant remanent polarization has been induced for the BiFeO<sub>3</sub> thin film with an optimum thickness, confirming that the BiFeO<sub>3</sub> thin film is a promising candidate material for the memory device.

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